

# Random Quantum Circuits are Unitary Polynomial-Designs

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# Outline

- **The problem**  
**Unitary  $t$ -designs**  
**Random Circuits**
- **Result**  
**Poly( $n$ ) Random Circuits are poly( $n$ )-designs**
- **Applications**  
**Fooling Small Sized Circuits**  
**Quick Equilibration by Unitary Dynamics**
- **Proof**  
**Connection to Spectral Gap of Local Hamiltonian**  
**A Lower Bound on the Spectral Gap**  
**Path Coupling for the Unitary Group**

# Haar Random Unitaries

For every integrable function in  $\mathbf{U}(d)$  and every  $V$  in  $\mathbf{U}(d)$

$$E_{U \sim \text{Haar}} f(U) = E_{U \sim \text{Haar}} f(VU)$$

# Applications of Haar Unitaries

(Hayden, Leung, Winter '04) Create **entangled states** with **extreme properties**

(Emerson et al '04) **Process tomography**

(Hayden et al '04) **Quantum data hiding** and **information locking**

(Sen '05) **State distinguishability**

(Abeyesinghe '06) Encode for **transmission** of quantum information through a quantum channel, **state merging**, **mother protocol**, ...

# The Price You Have to Pay...

To sample from the Haar measure with error  $\epsilon$  you need  $\exp(4^n \log(1/\epsilon))$  different unitaries

**Exponential** amount of random bits and quantum gates...

# Quantum Pseudo-Randomness

In many applications, we can replace a Haar random unitary by *pseudo-random* unitaries:

**This talk:** Quantum Unitary  $t$ -designs

**Def.** An ensemble of unitaries  $\{\mu(dU), U\}$  in  $\mathbf{U}(d)$  is an  $\varepsilon$ -approximate unitary  $t$ -design if for every monomial

$$M = U_{p1, q1} \dots U_{pt, qt} U_{r1, s1}^* \dots U_{rt, st}^*$$

$$|E_{\mu}(M(U)) - E_{\text{Haar}}(M(U))| \leq d^{-2t}\varepsilon$$

# Quantum Unitary Designs

Conjecture 1. There are **efficient**  $\varepsilon$ -approximate unitary  $t$ -designs  $\{\mu(dU), U\}$  in  $\mathbf{U}(2^n)$

Efficient means:

- unitaries created by  **$\text{poly}(n, t, \log(1/\varepsilon))$**  two-qubit gates
- $\mu(dU)$  can be sampled in  **$\text{poly}(n, t, \log(1/\varepsilon))$**  time.

# Quantum Unitary Designs

## Previous work:

(DiVincenzo, Leung, Terhal '02) Clifford group is an *exact* 2-design

(Dankert et al '06) Efficient construction of 2-design

(Ambainis and Emerson '07) Efficient construction of *state*  
poly(n)-design

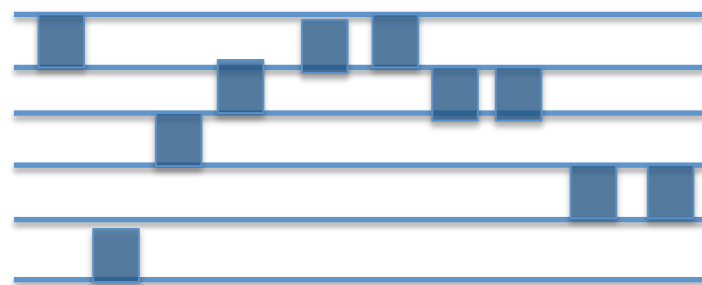
(Harrow and Low '08) Efficient construction of (n/log(n))-design

(Abeyesinghe '06) 2-designs are enough for decoupling

(Low '09) Other applications of t-design (mostly 2-designs)  
replacing Haar unitaries

# Random Quantum Circuits

**Local Random Circuit:** in each step an index  $i$  in  $\{1, \dots, n\}$  is chosen uniformly at random and a two-qubits Haar unitary is applied to qubits  $i$  e  $i+1$



**Random Walk** in  $\mathbf{U}(2^n)$

(Another example: **Kac's random walk** – toy model Boltzmann gas)

**Introduced** in (Hayden and Preskill '07) as a toy model for the dynamics of a black hole

# Random Quantum Circuits

## Previous work:

(Oliveira, Dalhsten, Plenio '07)  $O(n^3)$  random circuits are 2-designs

(Harrow, Low '08)  $O(n^2)$  random Circuits are 2-designs for every universal gate set

(Arnaud, Braun '08) numerical evidence that  $O(n \log(n))$  random circuits are unitary  $t$ -design

(Znidaric '08) connection with spectral gap of a mean-field Hamiltonian for 2-designs

(Brown, Viola '09) connection with spectral gap of Hamiltonian for  $t$ -designs

(B., Horodecki '10)  $O(n^2)$  local random circuits are 3-designs

# Random Quantum Circuits as $t$ -designs?

Conjecture 2. Random Circuits of size  $\text{poly}(n, \log(1/\epsilon))$  are an  $\epsilon$ -approximate unitary  $\text{poly}(n)$ -design

# Main Result

Conjecture 2. Random Circuits of size  $\text{poly}(n, \log(1/\epsilon))$  are an  $\epsilon$ -approximate unitary  $\text{poly}(n)$ -design

(B., Harrow, Horodecki '11) Local Random Circuits of size  $\tilde{O}(n^2 t^5 \log(1/\epsilon))$  are an  $\epsilon$ -approximate unitary  $t$ -design

# Data Hiding

## Computational Data Hiding:

“Most quantum states look maximally mixed for all polynomial sized circuits”

e.g. most quantum states are useless for measurement based quantum computation (Gross et al '08, Bremner et al '08)

Let  $QC(k)$  be the set of 2-outcome POVM  $\{A, I-A\}$  that can be implemented by a circuit with  $k$  gates

$$\Pr_{|\psi\rangle \sim \text{Haar}} \left( \max_{A \in QC(\text{poly}(n))} \left| \langle \psi | A | \psi \rangle - 2^{-n} \text{tr}(A) \right| \geq \varepsilon \right) \leq 2^{n \log(n)} 4^{-c\varepsilon^2 2^n}$$

# Data Hiding

## Computational Data Hiding:

“Most quantum states look maximally mixed for all polynomial sized circuits”

1. By **Levy's Lemma**, for every  $0 < A < I$ ,

$$\Pr_{|\psi\rangle \sim \text{Haar}} (|\langle \psi | A | \psi \rangle - 2^{-n} \text{tr}(A)| \geq \varepsilon) \leq e^{-c\varepsilon^2 2^n}$$

2. There is a eps-net of size  $< \exp(n^{\log(n)})$  for poly(n) implementable POVMs. By **union bound**

$$\Pr_{|\psi\rangle \sim \text{Haar}} \left( \max_{A \in \text{QC}(\text{poly}(n))} |\langle \psi | A | \psi \rangle - 2^{-n} \text{tr}(A)| \geq \varepsilon \right) \leq 2^{n^{\log(n)}} 4^{-c\varepsilon^2 2^n}$$

# Data Hiding

## Computational Data Hiding:

“Most quantum states look maximally mixed for all polynomial sized circuits”

1. By Levy's Lemma for every  $0 < A < I$ ,

$\Pr_{|\psi\rangle \sim \text{Haar}}$

But most states also require  $2^{O(n)}$  quantum gates to be approximately created...

2. There is a  $\epsilon$ -net of size  $\leq \exp(n \epsilon^{-2})$  for  $\text{poly}(n)$  implementable POVMs. By union bound

$$\Pr_{|\psi\rangle \sim \text{Haar}} \left( \max_{A \in \text{QC}(\text{poly}(n))} |\langle \psi | A | \psi \rangle - 2^{-n} \text{tr}(A)| \geq \epsilon \right) \leq 2^{n \log(n)} 4^{-c\epsilon^2 2^n}$$

# Efficient Data Hiding

**Corollary 1:** Most quantum states formed by  $O(n^k)$  circuits look maximally mixed for every circuit of size  $O(n^{(k+4)/6})$

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**Same idea** (small probability + eps-net), but replace Levy's lemma by a  $t$ -design bound from (Low '08):

$$\Pr_{U \sim \nu_{s,n}} \left( \left| \langle 0 | UAU | 0 \rangle - 2^{-n} \text{tr}(A) \right| \geq \delta \right) \leq \exp(O(t \log(1/\delta) - nt))$$

with  $t = s^{1/6} n^{-1/3}$  and  $\nu_{s,n}$  the measure on  $U(2^n)$  induced by  $s$  steps of the local random circuit model

$\epsilon$ -net of POVMs with  $r$  gates has size  $\exp(O(r(\log(n) + \log(1/\epsilon))))$

# Circuit Minimization Problem

**Goal:** Given a unitary, what is the minimum number of gates needed to approximate it to an error  $\epsilon$ ?

**Circuit Complexity:**

$$C_\epsilon(U) := \min\{k : \text{there exists } V \text{ with } k \text{ gates s.t. } \|V - U\| \leq \epsilon\}$$

**Question:** Lower bound to the circuit complexity?

**Corollary 2:** Most circuits of size  $O(n^k)$  have  $C_\epsilon(U) > O(n^{(k+4)/6})$

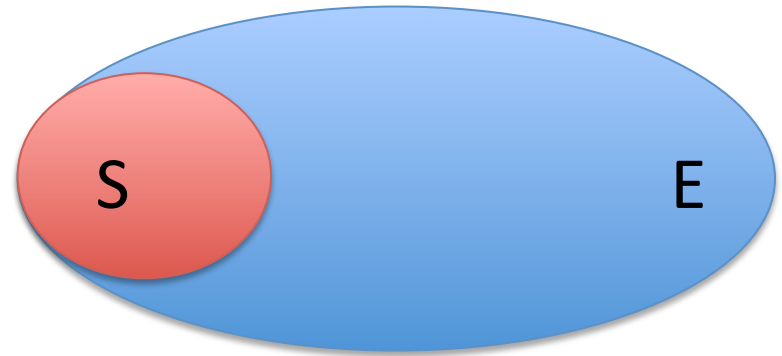
# Haar Randomness Not Needed

More generally,

Any quantum algorithm that has  $m$  uses of a Haar unitary and  $l$  gates and accepts, on average, with probability  $p$ , will accept with probability in  $(p - \epsilon, p + \epsilon)$  if we replace the Haar unitary by a random circuit of size  $\text{poly}(m, l, \log(1/\epsilon))$

# Equilibration by Unitary Dynamics

**Problem:** Let  $H_{SE}$  be a Hamiltonian of two quantum systems, S and E with  $|S| \ll |E|$



State at time  $t$ :

$$\rho_S(t) := \text{tr}_E \left( e^{-itH_{SE}} \rho_0 e^{itH_{SE}} \right)$$

On physical grounds we expect that for most times

$$\rho_S(t) \approx \rho_{\text{equi}}$$

This is true, in the limit of infinite times! (Linden et al '08)

# Fast Equilibration by Unitary Dynamics

How about the **time scale** of equilibration?  
For which  $T$  do we have

$$\frac{1}{T} \int_0^T \|\rho_S(t) - \rho_{equi}\| dt \leq \varepsilon \quad ?$$

(Linden et al '08) only gives the bound  $T \leq 1/(\text{min. energy gap})$

But we know equilibration is fast:  
coffee gets cold quickly, beer gets warm quickly ☹️

# Fast Equilibration by Unitary Dynamics

Toy model for equilibration: Let  $H_{SE} = UDU'$ , with  $U$  taken from the Haar measure in  $U(|S||E|)$  and  $D := \text{diag}(E_1, E_2, \dots)$ .

$$\frac{1}{T} \int_{U(|S||E|)} \text{tr} \left( \left( \rho_S(t) - \rho_{\text{equi}} \right)^2 \right) \mu_{\text{Haar}}(dU) \leq O \left( \frac{\left( \sum_k e^{i2tE_k} \right)^2}{|S||E|^2} + \dots \right)$$

(B., Ciwiklinski et al '11, Masanes et al '11, Vinayak, Znidaric '11)

Time of equilibration: Average energy gap:  $\frac{1}{|S|^2|E|^2} \sum_{j,l} (E_j - E_l)^{-1}$

For typical eigenvalue distribution goes with  $O(1/\log(|E|))$

# Fast Equilibration by Unitary Dynamics

Calculation only involves **4th** moments:

$$\frac{1}{T} \int_{U(|S||E|)} \text{tr} \left( \left( \text{tr}_E \left( U e^{-itD} U' \rho_0 U e^{itD} U' \right) - \rho_{\text{equi}} \right)^2 \right) \mu_{\text{Haar}}(dU)$$

Can replace Haar measure by an approximate unitary **4**-design

**Corollary 3.** For most Hamiltonians of the form  $UDU'$  with  $U$  a random quantum circuit of  $O(n^3)$  size, small subsystems equilibrate fast.

# Fast Equilibration vs Diagonalizing Complexity

Let  $H = UDU'$ , with  $D$  diagonal. Then we call  $C_\varepsilon(U)$  the diagonalizing complexity of  $U$ .

**Corollary 4.** For most Hamiltonians with  $O(n^3)$  diagonalizing complexity, small subsystems equilibrate fast.

Connection suggested in (Masanes, Roncaglia, Acin '11)

**In contrast:** Hamiltonians with  $O(n)$  diagonalizing complexity  
Do not equilibrate in general

**Open question:** Can we prove something for the more interesting case of **few-body** Hamiltonians?

# Proof of Main Result

1. **Another characterization** of unitary t-designs
2. Mapping the problem to bounding spectral gap of a **Local Hamiltonian**
3. Technique for **bounding spectral gap**

“It suffices to get a *exponential small* bound on the convergence rate”

4. **Path Coupling** applied to the unitary group

# t-Copy Tensor Product Quantum Expanders

An ensemble of unitaries  $\{\mu(dU), U\}$  is an  $(t, 1-\varepsilon)$  tensor product expander if

$$\left\| \int \mu(dU) U^{\otimes t} \otimes \bar{U}^{\otimes t} - \int \mu_H(dU) U^{\otimes t} \otimes \bar{U}^{\otimes t} \right\|_{\infty} \leq \varepsilon$$

Obs: Implies it is a  $d^{2t}\varepsilon$ -approximate unitary  $t$ -design

# Relating to Spectral Gap

$\mu_n$ : measure on  $U(2^n)$  induced by one step of the local random circuit model

$(\mu_n)^{*k}$ :  $k$ -fold convolution of  $\mu_n$  (measure induced by  $k$  steps of the local random circuit model)

We show:

$$\left\| \int \mu_n^{*k}(dU) U^{\otimes t} \otimes \bar{U}^{\otimes t} - \int \mu_H(dU) U^{\otimes t} \otimes \bar{U}^{\otimes t} \right\|_{\infty} = \lambda_2 \left( \int \mu_n(dU) U^{\otimes t} \otimes \bar{U}^{\otimes t} \right)^k$$

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It suffices to prove an upper bound on  $\lambda_2$  of the form

$1 - \Omega(t^{-4}n^{-1})$  since  $(1 - \Omega(t^{-4}n^{-2}))^k \leq 2^{-2nt}\epsilon$  for  $k = O(n^2t^5 \log(1/\epsilon))$

# Relating to Spectral Gap

But 
$$\mu_n = \frac{1}{n} \sum_{i=1}^n \mu_{Haar}(i, i+1)$$

So 
$$\lambda_2 \left( \int \mu_n(dU) U^{\otimes t} \otimes \bar{U}^{\otimes t} \right) = 1 - \frac{\Delta(H_{n,t})}{n}$$

with 
$$H_{n,t} := \sum_{i=1}^n h_{i,i+1} \quad h_{i,i+1} := I - \int_{U(4)} U_{i,i+1}^{\otimes t} \otimes \bar{U}_{i,i+1}^{\otimes t} \mu_H(dU)$$

and  $\Delta(H_{n,t})$  the spectral gap of the local Hamiltonian  $H_{n,t}$



# Relating to Spectral Gap

But 
$$\mu_n = \frac{1}{n} \sum_{i=1}^n \mu_{Haar}(i, i+1)$$

So 
$$\lambda_2 \left( \int \mu_n(dU) U^{\otimes t} \otimes \bar{U}^{\otimes t} \right) = 1 - \frac{\Delta(H_{n,t})}{n}$$

wit  $J_{i,i+1}^{\otimes t} \otimes \bar{U}_{i,i+1}^{\otimes t} \mu_H(dU)$   
 and  $H_{n,t}$  Hamiltonian  $H_{n,t}$

Want to lower bound spectral gap by  $O(t^{-4})$



# Structure of $H_{n,t}$

i. The minimum eigenvalue of  $H_{n,t}$  is zero and the **zero eigenspace** is

$$G_{n,t} := \text{span} \left\{ |\psi_\pi\rangle^{\otimes n}, |\psi_\pi\rangle := (I \otimes V(\pi)) |\Phi(2^t)\rangle : \pi \in S_t \right\}$$

ii. **Approximate orthogonality** of permutation matrices:

$$\sum_{\pi \in S_t} \left| \langle \psi_\sigma | \psi_\pi \rangle \right|^n \leq 1 + \frac{2t^2}{2^n}, \quad \left\| \sum_{\pi \in S_t} (|\psi_\pi\rangle \langle \psi_\pi|)^{\otimes n} - G_{n,t} \right\|_\infty \leq \frac{2t^2}{d^n}$$

# Structure of $H_{n,t}$

Follows from

$$[X, U^{\otimes t}] = 0 \Leftrightarrow X = V(\pi)$$

i. The minimum eigenvalue  
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# Structure of $H_{n,t}$

Follows from

$$[X, U^{\otimes t}] = 0 \Leftrightarrow X = V(\pi)$$

i. The minimum eigenvalue

Follows from

$$P_{sym} = \frac{1}{t!} \sum_{\pi \in S_t} V(\pi), \quad \{ |\psi_\pi\rangle := (I \otimes V(\pi)) |\Phi(2^t)\rangle : \pi \in S_t \}$$

ii. Approximate orthogonality of permutation matrices:

$$\sum_{\pi \in S_t} |\langle \psi_\sigma | \psi_\pi \rangle|^n \leq 1 + \frac{2t^2}{2^n}, \quad \left\| \sum_{\pi \in S_t} (|\psi_\pi\rangle \langle \psi_\pi|)^{\otimes n} - G_{n,t} \right\|_\infty \leq \frac{2t^2}{d^n}$$

# Lower Bounding $\Delta(H_{n,t})$

We prove: 
$$\Delta(H_{n,t}) \geq \frac{\Delta(H_{2\log(t),t})}{8\log(t)}$$

Follows from [structure of  \$H\_{n,t}\$](#)  and

(Nachtergaele '96) Suppose there exists  $l, n_l, \varepsilon_l$  such that

$$\text{for all } n_l < m < n-1 \quad \left\| (I_{A_1} \otimes G_{A_2 B})(G_{A_1 A_2} \otimes I_B) - G_{A_1 A_2 B} \right\|_{\infty} \leq \varepsilon_l$$

with  $A_1 := [1, m-l-1]$ ,  $A_2 := [m-l, m-1]$ ,  $B := m$  and  $\varepsilon_l < l^{-1/2}$ . Then

$$\Delta(H_{[1,n]}) \geq \Delta(H_{[1,l]}) \left( \frac{(1 - \varepsilon_l \sqrt{l})}{l-1} \right)$$

# Lower Bounding $\Delta(H_{n,t})$

We prove: 
$$\Delta(H_{n,t}) \geq \frac{\Delta(H_{2\log(t),t})}{8\log(t)}$$

Follow (Nac) Want to lower bound by  $O(t^{-4})$ , an exponential small bound in the size of the chain (i.e. in  $2\log(t)$ )

for all  $n_l < m < n-1$  
$$\left\| (I_{A_1} \otimes G_{A_2 B})(G_{A_1 A_2} \otimes I_B) - G_{A_1 A_2 B} \right\|_{\infty} \leq \varepsilon_l$$

with  $A_1 := [1, m-l-1]$ ,  $A_2 := [m-l, m-1]$ ,  $B := m$  and  $\varepsilon_l < l^{-1/2}$ . Then

$$\Delta(H_{[1,n]}) \geq \Delta(H_{[1,l]}) \left( \frac{(1 - \varepsilon_l \sqrt{l})}{l-1} \right)$$

# Exponentially Small Bound to Spectral Gap

Follows from two relations:

$$1. \left(1 - \frac{\Delta(H_{n,t})}{n}\right)^k \leq 2tW\left(\left(\mu_n\right)^{*k}, \mu_{Haar}\right)$$

Wasserstein distance:

$$W(\nu_1, \nu_2) := \sup\left\{\int f(u)\nu_1(du) - \int f(u)\nu_2(du) : f \text{ is } 1\text{-Lipschitz}\right\}$$

$$2. W\left(\left(\mu_n\right)^{*{(n-1)k}}, \mu_{Haar}\right) \leq 2^{n/2} (1 - 2^{-5n})^{\frac{k}{n-1}}$$

# Bounding Convergence with Path Coupling

**Key result** to 2<sup>nd</sup> relation: Extension to the unitary group of Bubley and Dyer **path coupling**

Let  $W_p(\nu_1, \nu_2) := \inf \left\{ E[d(X, Y)^p]^{1/p} : (X, Y) \text{ couples } (\nu_1, \nu_2) \right\}$

(Oliveira '07) Let  $\nu$  be a measure in  $U(d)$  s.t.

$$\lim_{\varepsilon \rightarrow 0} \sup_{U_1, U_2 \in U(d)} \left\{ \frac{W_2(\nu * \delta_{U_1}, \nu * \delta_{U_2})}{\|U_1 - U_2\|_2} : \|U_1 - U_2\|_2 \leq \varepsilon \right\} \leq \eta$$

Then  $W_2(\nu * \nu_1, \nu * \nu_2) \leq \eta W_2(\nu_1, \nu_2)$

# Bounding Convergence with Path Coupling

Must consider coupling in  $n$  steps of the walk to get non trivial contraction (see paper for details)

(Oliveira '07) Let  $\nu$  be a measure in  $U(d)$  s.t.

$$\lim_{\varepsilon \rightarrow 0} \sup_{U_1, U_2 \in U(d)} \left\{ \frac{W_2(\nu * \delta_{U_1}, \nu * \delta_{U_2})}{\|U_1 - U_2\|_2} : \|U_1 - U_2\|_2 \leq \varepsilon \right\} \leq \eta$$

Then  $W_2(\nu * \nu_1, \nu * \nu_2) \leq \eta W_2(\nu_1, \nu_2)$

# Summary

- $\tilde{O}(n^2 t^5 \log(1/\epsilon))$  local random circuits are  $\epsilon$ -approximate unitary  $t$ -designs
- Most states of size  $n^k$  is indistinguishable from maximally mixed by all circuits of size  $n^{(k+4)/6}$
- Proof is based on
  - (i) connection to spectral gap local Hamiltonian
  - (ii) approximate orthogonality of permutation matrices
  - (iii) path coupling for the unitary group
- Another application to fast equilibration of quantum systems by unitary dynamics with an environment

# Open Questions

- Is  $\tilde{O}(n^2 t^5 \log(1/\epsilon))$  tight?
- Can we prove that **constant depth** random circuits are approximate unitary  $t$ -designs?  
(we can show they form a  $t$ -tensor product expander;  
proof uses the **detectability lemma** of Aharonov et al)

Would have **applications** to:

- (i) fast equilibration of generic **few-body** Hamiltonians
- (ii) creation of topological order by short circuits  
(counterpart to the no-go result of Bravyi, Hastings, Verstraete for short *local* circuits)